Modenizing Quantum Annealing using Local Search

PC 2016 Manchester See full paper at: arXiv:1606.06833

Nick Chancellor

July 14, 2016



イロト 不得 トイヨト イヨト

Outline

- 1. Energy Computing and the Ising Model
- 2. Quantum Annealing and Simulated Annealing
 - Better Classical Algorithms: Parallel Tempering and Population Annealing
 - Hybrid Computing: Gaining the Advantages of Advanced Algorithms
- 3. Search Range
 - Controlling Range
 - Parallel tempering and population annealing analogues

- Controlling Problem mis-specification
- 4. A few slides about implementation
- 5. Thermal Sampling

Problem Statement: Ising Spin Glass Hamiltonian

$$H_{ISG} = \sum_{i} h_i \sigma_i^z + \sum_{ij} J_{ij} \sigma_i^z \sigma_j^z$$

Goal is to find ground/low energy states

- 'Universal' in the sense that any classical Hamiltonian can be mapped to it De las Cuevas, Cubitt Science 351 6278
- Thermal/quantum distributions also useful for inference and machine learning tasks ex. Amin et. al. arXiv:1601.02036, Chancellor et. al. Scientific Reports 6, 22318 ...



Simulated Annealing (classical)

Updates drive toward thermal distribution with temperature T if they obey detailed balance

$$P(S(1) \to S(2)) = \exp(\frac{(E(1) - E(2))}{T})P(S(2) \to S(1))$$

Start at high T and lower over time



▲□▶ ▲圖▶ ★ 国▶ ★ 国▶ - 国 - のへで

Quantum Annealing (QA)



Add non-commuting transverse field terms

$$H(s) = -A(s)\sum_{i}\sigma_{i}^{x} + B(s)H_{ISG}$$

start at $\frac{A(s=0)}{B(s=0)} \gg 1$, go to $\frac{B(s=1)}{A(s=1)} \gg 1$ Quantum fluctuations + low temperature bath cause tunneling toward low energy states

Beyond Simulated Annealing (classical)

Parallel Tempering

Multiple replicas at different temperatures



 Swap replicas by rules which obey detailed balance

$$P_{swap}(i,j) = \\ \min\left[1, \exp\left(\left(\frac{1}{T(i)} - \frac{1}{T(j)}\right)(E_i - E_j)\right)\right]$$

Population Annealing

Anneal multiple replicas



 Probabilistically remove poorly performing replicas and copy those which perform well
Rules preserve average population and obey detailed balance

•
$$\bar{N}(E) = \frac{1}{Q} \exp\left(\left(\frac{1}{T_{old}} - \frac{1}{T_{new}}\right)E\right)$$

Can hybrid strategies combine these with calls to an annealer? Can these strategies be used directly by a quantum annealer?

Difficulties in building new annealer strategies

- \blacktriangleright No cloning theorem \rightarrow cannot copy quantum states
- Measurements (ex. energy) disturb state of system and likely experimentally difficult
- Usual QA is global search, no way of inserting information from previous runs

Solution \rightarrow use annealer subroutine which starts and ends at s = 1 (recall $\frac{B(s=1)}{A(s=1)} \gg 1$) with programmed initial state



Hybrid computing using local search

Potential Strategies

1. Quantum and classical algorithms used together

- Classical input and output means that annealer can be used alongside any classical algorithm
- Google have started looking into these ideas, see Hartmut Neven talk at AQC 2016¹
- 2. Multiple local quantum searches controlled by classical algorithm
 - Analogues to parallel tempering and population annealing which use annealer only
 - Will return to this later

¹Should be uploaded soon and viewable at: https://aqc2016.eventfarm.com \circ and \circ

Cartoon example: energy landscape with rough and smooth features



- a) QA gets stuck in broad local minima and cannot tunnel to correct minima
- b) Classical algorithms can easily explore the broad features, while the annealer can explore the rough ones
- c) Even random initialization can improve solution probabilities, may hit rough region by chance

Range of local search

Define search range in terms of typical Hamming distance $\mathfrak{h}(s')$ from starting state

 $\mathfrak{h}(s')$ will increase monotonically with decreasing s' but...

- Not easy to theoretically predict
- Will depend on both problem and starting state



▲ロト ▲帰 ト ▲ ヨ ト ▲ ヨ ト ・ ヨ ・ の Q ()

Choosing the range of the search

Options:

- 1. Choose heuristically: use different ranges and take best and take best, typical $\mathfrak{h}(s')$ for problem type etc...
- 2. Measure search range and use bisection to get to desired range
- 3. Define effective temperature and construct analogues of known classical algorithms



Effective Temperature

1. Single qubit Hamiltonian with transverse and longitudinal components

$$H_1(s') = -A(s')\,\sigma^x + B(s')\,\sigma^z$$

2. Diagonalize 2x2 matrix by hand to get occupation ratio

$$\frac{\psi(1)}{\psi(2)} = \frac{\sqrt{A(s')^2 + B(s')^2}}{A(s')} + \frac{B(s')}{A(s')}$$

3. Invert Boltzmann equation to get effective temperature

$$T_{eff}(s') = 2 \left[\ln \left(\left| \frac{\psi(1)}{\psi(2)} \right|^2 \right) \right]^{-1}$$

Parallel tempering and population annealing analogues

- 1. Replace metropolis updates with annealing runs consisting of calls to annealer
- 2. Define 'energy' and 'state' as the lowest energy solution found in an annealing run and the corresponding classical state

- 3. Replace $T \rightarrow T_{eff}$
- 4. Apply replica swapping/copying/deleting rules as usual



Problem mis-specification

- Error in each energy proportional to $\sqrt{N_{qubit}}$
- Only energy differences within search matter
- Energy difference proportional to square root of Hamming distance
- : relevant error proportional to $\sqrt{\mathfrak{h}(s')}$ not $\sqrt{N_{qubit}}$ ²



 $^{^{2}}$ Up to details about shape of the explored subspace; see arXiv:1606.06833

Implementation background: flux qubit circuit

Compound Compound Josephson Junction device like those used in D-Wave systems



- ϕ is relevant variable
- ϕ_c Controls effective transverse field
- ϕ^x Acts as external bias
- Couple inductively between loops

Implementation Background: Single qubit potential



- Barrier width controlled by ϕ_c
- Energy difference between wells controlled by ϕ^x
- Quantum tunneling suppressed exponentially late in the anneal

・ロト ・ 雪 ト ・ ヨ ト

э

High barrier also blocks classical transitions

Runback protocol



1. Anneal forward using standard annealing protocol and trivial Hamiltonian to initialize state

- 2. Reprogram problem Hamiltonian to target problem
 - State protected due to high energy barrier
- 3. Anneal back to point s'
- 4. Possibly wait a period au and anneal back to s=1

One slide on sampling

Annealer calls will not obey detailed balance, but...

- Somtimes quantum distributions can act as an effective proxy for thermal distributions Otsubo et. al. Phys. Rev. E 86, 051138
- Quantum fluctuations may aid in machine learning tasks Amin et. al. arXiv:1601.02036

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

 Relative weights of local minima can be calculated though post processing: numerically calculate free energy with classical Monte Carlo

Acknowedgements

- Thanks to Viv Kendon for multiple critical readings of the paper
- Work supported by EPSRC
- You \rightarrow thanks for listening

You are encouraged to read the full paper: arXiv:1606.06833

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <