

max-SAT, multi-body frustration, and multi-body sampling on a two local Ising system

AQC 2016 Los Angeles

Nick Chancellor

June 29, 2016



Acknowledgements

Collaborators:

Stefan Zohren (Oxford)
Paul Warburton (UCL)
Simon Benjamin (Oxford)
Stephen Roberts (Oxford)

Supervisor:

Viv Kendon (Durham)

Funding:

NC: UK Engineering and Physical Sciences Research Council
(EPSRC)

SZ: Lockheed Martin, Nokia, and EPSRC

PW: Lockheed Martin and EPSRC

SB: EPSRC

Outline

1. SAT, max-SAT, frustration, non-linear constraints, and sampling
2. Spectral Mapping with Ising Spins
 - ▶ Arbitrary max-SAT on the chimera graph
 - ▶ Aside: Another way to do parity checks
 - ▶ A better mapping
3. Applications
 - ▶ Classical communications: turbo code decoder
 - ▶ Particle simulation

Penalty formalism for SAT

Hamiltonian to implement clauses:

$$H(a) = \sum_i Pen(\{a_i^{(l)}\})$$

$$a_i^{(l)} \in \{a_1, a_2, a_2, \dots\} \cup \{\neg a_1, \neg a_2, \neg a_3, \dots\}$$

$$Pen(\{a^{(l)}\}) \begin{cases} \geq g & a_i^{(l)} = 0, \forall i \\ = 0 & \text{otherwise.} \end{cases}$$

- ▶ Violating any clause gives an energy of at least g
- ▶ zero energy state satisfies all clauses
- ▶ Lowest energy state not meaningful if cannot satisfy simultaneously

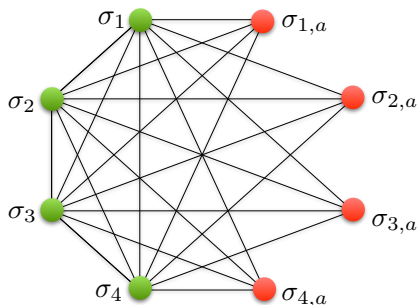
Alternate approach: penalize every clause violation equally

$$Spec(\{a^{(l)}\}) = \begin{cases} g & a_i^{(l)} = 0, \forall i \\ 0 & \text{otherwise.} \end{cases}$$

$$H(a) = \sum_i Spec(\{a_i^{(l)}\}) = N(a)g$$

- ▶ $N(a)$ is the number of clauses violated by bitstring a
- ▶ minimum energy bitstring satisfies the maximum number of clauses
- ▶ Boltzmann distribution can be used for maximum entropy inference → will come back to this

Making clauses using the Ising model ([arXiv:1604.00651](https://arxiv.org/abs/1604.00651))



- ▶ couple all logical bits to each other and to all ancillas
- ▶ bias ancillas so that ancilla 1 will flip if 1 or more logical bits are up, 2 if 2 etc...
- ▶ ancilla couplings and couplings between bits cancel, fields on ancillas can assign arbitrary penalties
- ▶ single bit sector corresponds to clause on previous slide, can also create parity checks (\oplus clauses) [arXiv:1603.09521](https://arxiv.org/abs/1603.09521)

simple example

Logical bit values	Ancilla values	E
1111	0000	0
0111, 1011, 1101, 1110	0001	0
0011, 0101, 0110, 1001, 1010, 1100	0011	0
1000, 0100, 0010, 0001	0111	0
0000	1111	g

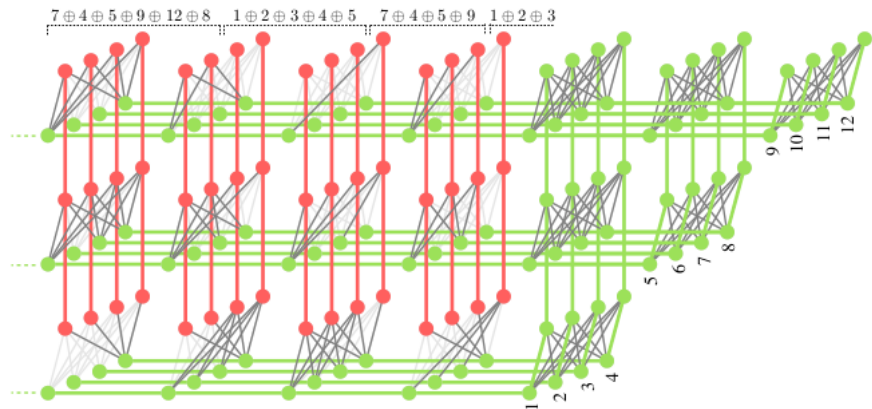
- ▶ Example in which applying a field on the last ancilla applies an OR clause $q_1 \vee q_2 \vee q_3 \vee q_4$
- ▶ Applying to first ancilla will instead apply AND $q_1 \wedge q_2 \wedge q_3 \wedge q_4$
- ▶ Alternating ancilla fields can apply XOR $q_1 \oplus q_2 \oplus q_3 \oplus q_4$
- ▶ see: [arXiv:1604.00651](https://arxiv.org/abs/1604.00651) for Hamiltonian details

Quick Aside: Non-linear Constraints

Logical bit values	Ancilla values	E
1111	0000	E_4
0111, 1011, 1101, 1110	0001	E_3
0011, 0101, 0110, 1001, 1010, 1100	0011	E_2
1000, 0100, 0010, 0001	0111	E_1
0000	1111	E_0

- ▶ Energies can be assigned arbitrarily: can implement non-linear constraints on the number of ones

max-SAT on the chimera



- ▶ Use Choi complete graph minor embedding plus rows of ancillas to implement clauses [arXiv:1604.00651](https://arxiv.org/abs/1604.00651)

Drawbacks of this technique

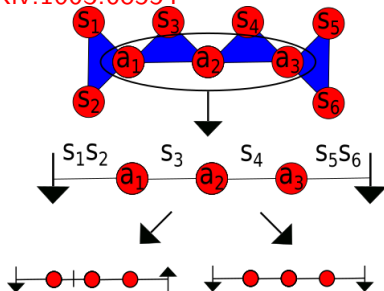
Conceptually nice to demonstrate direct mapping but....

- ▶ high connectivity makes mapping inefficient
- ▶ some clauses, ex. $(a_1 \vee a_2) \wedge (a_3 \vee a_4) \wedge \dots$ require an exponentially growing number of ancillas

I will show how both these problems can be addressed after a brief (1-slide) aside

one slide aside: another way of implementing parity checks

Proposed by other authors in [arXiv:1604.02359](https://arxiv.org/abs/1604.02359) , alternate formulation in [arXiv:1603.08554](https://arxiv.org/abs/1603.08554)

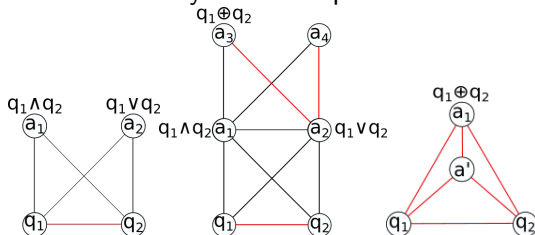


Triangles \rightarrow effective 3 body coupling gadgets, $s \rightarrow$ logical bit, $s_a \rightarrow$ ancillas

- ▶ parity of $\{s\}$ is even \rightarrow no domain wall (unique), odd \rightarrow domain wall (degenerate)
- ▶ degeneracy can be removed by weakening any of the 3 local gadgets, \rightarrow necessary for thermal sampling

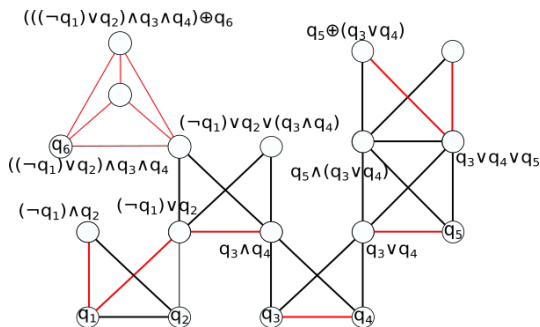
Ancilla representations for AND (\wedge), OR (\vee), and XOR (\oplus) on 2 bits

Using the methods of [arXiv:1604.00651](https://arxiv.org/abs/1604.00651), single 'indicator' bit corresponds to result of any of these operations



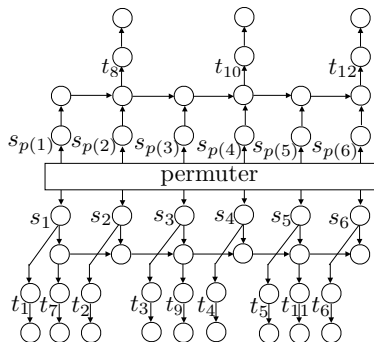
- ▶ \wedge and \vee 'mark' single states, so single bits act as 'indicators' for these clauses
- ▶ second copy of gadget Hamiltonian needed to create single bit corresponding to \oplus ,
- ▶ alternatively 3 bit gadget acts as indicator for \oplus

Chaining 2 bit Hamiltonians to make more complex clauses



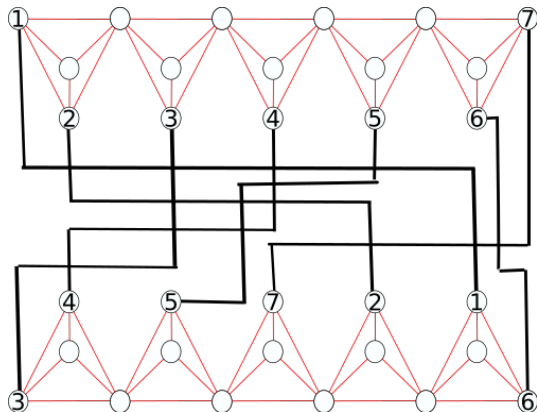
- ▶ Apply gadget to 'indicator' ancillas to create more complex clauses
- ▶ Fields on indicators enforce clauses
- ▶ **Any** clause which can be efficiently written using $\wedge, \vee, \oplus, \neg$, and parenthesis can be implemented efficiently
- ▶ Does not require full connectivity

Example: Classical Decoding, Turbo code implementation



- ▶ Interleaved convolution code: apply strings of \oplus of even length on bits $(s_1 \oplus s_2), (s_1 \oplus s_2 \oplus s_3 \oplus s_4) \dots$
- ▶ odd length strings of \oplus applied to random permutation $(s_{p(1)}), (s_{p(1)} \oplus s_{p(2)} \oplus s_{p(3)}) \dots$
- ▶ approaches Shannon limit for large block length

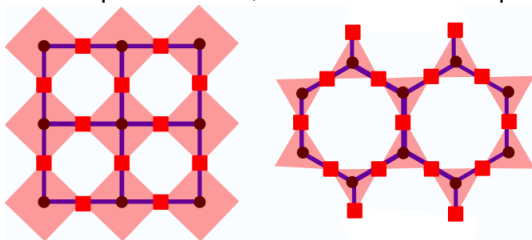
Turbo code decoder using minor embedding



- ▶ Black lines → embedding bonds
- ▶ red lines → ancilla couplings
- ▶ note similarity to ideas presented in [arXiv:1603.08554v2](https://arxiv.org/abs/1603.08554v2)

Application: Particle Simulation

Higher dimensional analogue of 1-D domain walls using high locality operators, red squares \rightarrow spins, pink polyhedra \rightarrow couplers, circles \rightarrow particle sites, lines \rightarrow allowed hops



$$H_{particle} = -\frac{m}{2} \sum_i \prod_{j \in d_i} \sigma_j^z - \tau \sum_i \sigma_i^z - \Delta \sum_i \sigma_i^x$$

Can this Hamiltonian be realized perturbatively using our gadget?

Realizing particle simulation Hamiltonian perturbatively

Using our gadget for each coupler

What works:

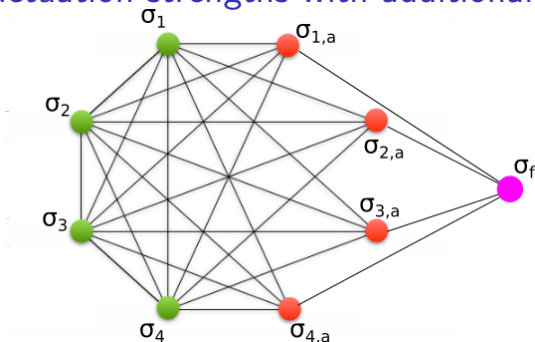
- ▶ All logical bit flips occur at 3rd order in perturbation theory (flip 2 ancillas plus bit)
- ▶ All transitions go through intermediate states with the same energy
- ▶ \therefore bit flips (σ_i^x) realized perturbatively with no modifications

What Doesn't:

- ▶ Second order fluctuations are not the same for every logical state \rightarrow some energetically favored over others

I will show how to rectify this on next slide

Matching fluctuation strengths with additional ancillas



'fluctuation control' ancillas (f), coupled to ancillas

$$H_{fc} = \sum_i (J_{i,f}(-\sigma_{i,a}^z + \sigma_{i,a}^z \sigma_f^z) + |J_{i,f}| \sigma_f^z) + h_f^0 \sigma_f^z$$

- ▶ $h_f^0 > 0$ no effect on Ising coupling
- ▶ $\{J_{i,f}\}$ tuned to control fluctuations
- ▶ more than 1 such ancilla can be added

Conclusions

- ▶ Alternate method for mapping problems based on max-SAT formalism
- ▶ Individual elements claimed together to realize any clause which can be written efficiently
- ▶ Example applications:
 - ▶ Classical message decoding: Turbo Code
 - ▶ Particle simulation: realize Hamiltonian perturbatively

Quick Plug for Some of my Other Work

Modernizing Quantum annealing using Local Search:
[arXiv:1606.06833](https://arxiv.org/abs/1606.06833)