

# Domain Walls in 1D: Simple Experiments to Answer Important Questions

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# Presentation Structure

1. Order-by-disorder effects on a forced domain wall
  - 1.1 “Particle In a Box” experiment
  - 1.2 Cause of the U-shape
  - 1.3 Using this to measure control errors
2. Domain wall tunneling
  - 2.1 Tunneling experiment
  - 2.2 Other uses of protocol

## Part 1.1: “Particle In a Box” experiment

Relatively simple spin chain Hamiltonian:



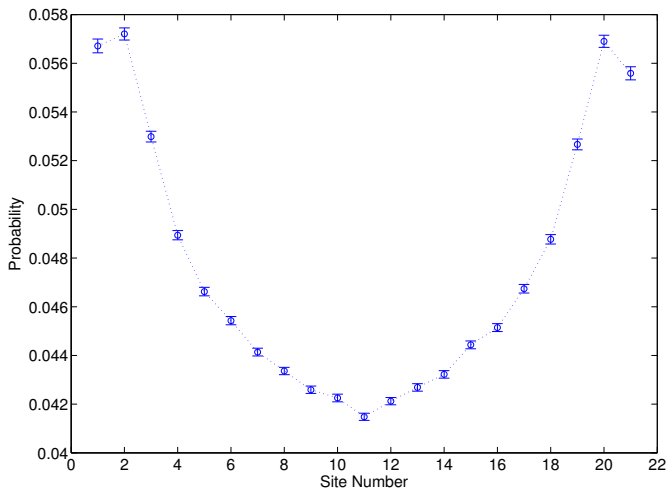
$$H(t) = B(t) \left( \sum_{i=1}^{N-1} -\sigma_i^z \sigma_{i+1}^z + 2(\sigma_1^z - \sigma_N^z) \right) - A(t) \sum_{i=1}^N \sigma_i^x$$

low energy sector for  $A \ll B$  has exactly one domain wall

Note, experiments have been done before with this Hamiltonian:

Johnson, M. W. et. al. Quantum annealing with manufactured spins. *Nature* **473**, 194-198 (12 May 2011).

# An unexpected result! U shaped domain wall distribution

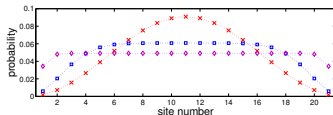


?!

# What could cause this? Usual suspects

## 1. Effect from transverse field -No!

- ▶ We expect the opposite shape from transverse field



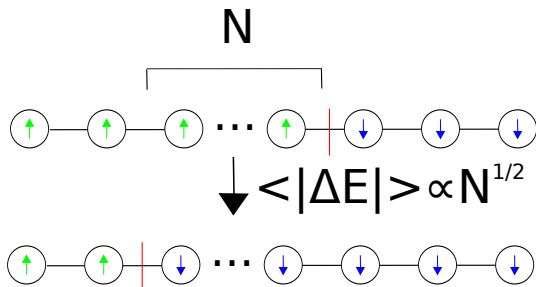
## 2. Annihilation -No!

- ▶ Can lead to a U-shape, but erased by diffusion
- ▶ Calculations and simulations predict equilibrium even for fastest anneal
- ▶ U-shape seen even for long annealing time

## 3. Control errors -How?

- ▶ Gauge and spacial averaging means that mean error will be the same everywhere on chain.
- ▶ Gauge averaging will make errors (approximately) uncorrelated.
- ▶ Can uncorrelated zero-mean errors do anything interesting?

## Ising domain wall subject to uncorrelated random fields



- ▶ Typical energy difference between states depends on number of spins which are different
- ▶ Pair of states with domain wall further away tends to have bigger energy difference

Domain wall energies are correlated even if fields are not!

## These correlations lead to a U-shape

- ▶ Mathematically: Taylor expand domain wall probability

$$P_n = \overline{e^{-\beta E_n}} = \left[ 1 + \sum_{m \neq n} e^{-\beta (E_m - E_n)} \right]^{-1} \text{ in } \beta \text{ to second order.}$$

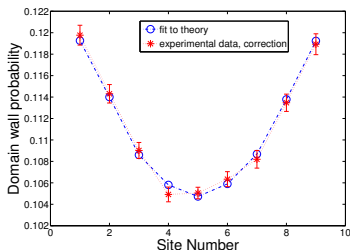
- ▶ Intuitively:

- ▶ Domain walls on nearby couplers tend to have similar energy
- ▶ Couplers on the end have fewer neighbours  $\rightarrow$  fewer domain wall locations with similar energy
- ▶  $\therefore Z = \sum_n \exp(-\beta E_n)$  tends to be smaller in cases where the domain wall has a high probability of being near the end



## Measuring control errors with U-shape

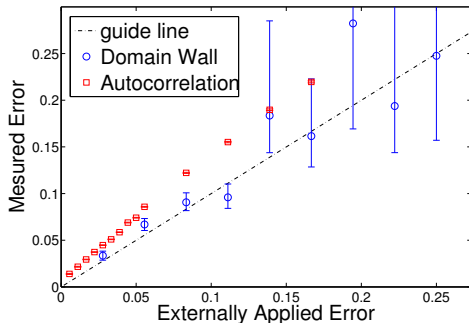
- ▶ For a chain of  $N$  spins,  $N - 1$  single domain wall states: Boltzmann distributions constructed exactly  $\rightarrow$  sample with noise of various strengths  $\rightarrow$  fit experiment to sampled



- ▶ Coupler control errors don't cause  $\{E\}$  to be correlated  $\rightarrow$  no effect on U
- ▶ Upper frequency limit related to equilibration timescale of the system and **not** to the annealing time

## How does this compare with other ways of measuring?

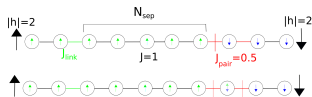
- ▶ Single qubit autocorrelation gives larger measurement of control error
- ▶ Also measures artificially introduced error more strongly



- ▶ Why the difference? one possibility is different freeze times

## Part II: Domain wall tunneling

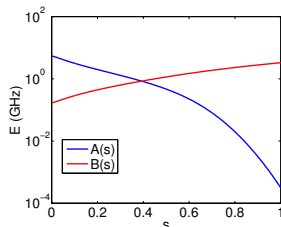
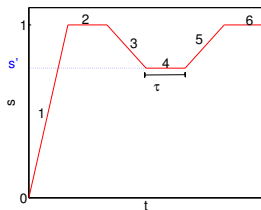
- ▶ Experiments which I will be performing over the summer at D-Wave headquarters in Burnaby B.C.
- ▶ Pattern Hamiltonian from part I with weak links to control domain wall resonances (more details later)



- ▶ More sophisticated annealing protocol allows initial state to be programmed
  - ▶ Extension of techniques used in Lanting, T. et. al. Entanglement in a Quantum Annealing Processor Phys. Rev. X 4 021041 (2014)

# Our experiment vs. typical operation

- ▶ Typical annealing run:  $h_i$  and  $J_{ij}$  programmed at  $s = 0 \rightarrow$  annealed linearly (in  $s$ ) to  $s = 1 \rightarrow$  final state read out
- ▶ Our protocol:
  1. Program strong  $h_i$  and no  $J_{ij}$  to set initial state, anneal  $s = 0 \rightarrow 1$
  2. At  $s = 1$  reprogram  $h_i$  and  $J_{ij}$  to target Hamiltonian
  3. Anneal back to  $0 < s' < 1$
  4. (optional) Wait for time  $\tau$
  5. Forward to  $s = 1$
  6. Measure



# Why would anyone want to do this?

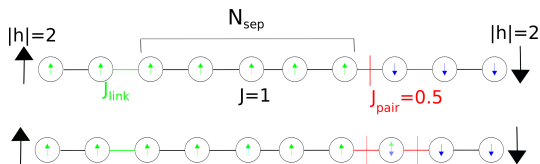
## Current project:

- ▶ Preparing in an excited state allows for simple tunneling experiments
- ▶ Probe “freezing” transition
- ▶ Directly access information about control error earlier in anneal

## Future projects:

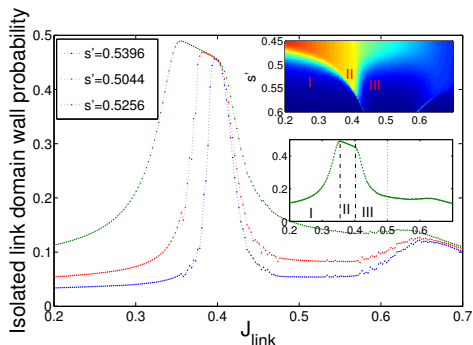
- ▶ Local search can be integrated with classical solvers in ways global search cannot
- ▶ Simulate diffusion of topological defects

# Domain wall tunneling with controlled resonance



- ▶ Place forced domain wall on pair of weak links
- ▶ Another weak link  $N_{sep}$  qubits away
- ▶ Infinitesimal transverse field causes qubit next to weak link to hybridize and reduce energy
- ▶ Pair and single link DW positions come into resonance at a value of  $s$  which depends on  $J_{link}$

# Open quantum system simulations predict 3 regimes

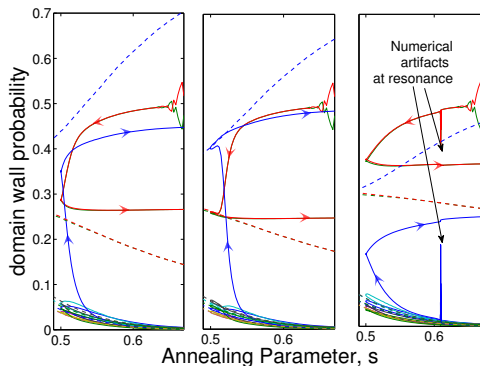


I Resonance limited: link and pair never get close enough to resonance for domain wall to tunnel

II Equilibration: resonance achieved at moderate transverse field; domain wall probabilities determined by detailed balance

III Freezing limited: resonance achieved at a point when transverse field too weak to facilitate tunneling




## Are these explanations correct?



- ▶ Red and Green lines are weak link pair, Blue line is isolated link
- ▶ Arrows indicate time direction
- ▶ Dashed lines are thermal equilibrium



# References

-  T. Lanting et. al. Entanglement in a Quantum Annealing Processor Phys. Rev. X 4 021041 doi: 10.1103/PhysRevX.4.021041 (2014)
-  R. Harris et. al. Experimental demonstration of a robust and scalable flux qubit PHYSICAL REVIEW B 81, 134510 (2010)
-  I Hen et. al. Probing for quantum speedup in spin glass problems with planted solutions arXiv:1502.01663 (2015)